

Supermatrix Models*



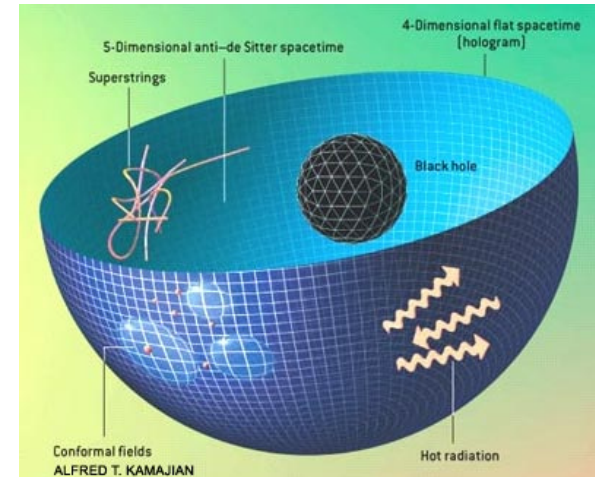
Robbert Dijkgraaf
Institute for Advanced Study

Walter Burke Institute for Theoretical Physics Inaugural Symposium
Caltech, Feb 24, 1015

* *Work with Cumrun Vafa and Ben Heidenreich, in progress*

Holography/Emergent Geometry

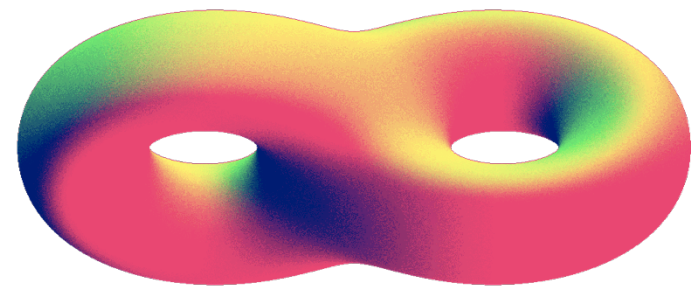
Quantum Theory
in d dimensions



$d+1$ Space-Time Gravity + QFT

Matrix Models ($d=0$)

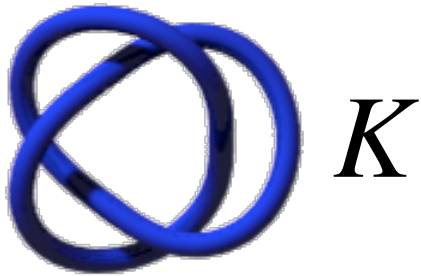
$$\int_{N \times N} d\Phi \cdot e^{\text{Tr } W(\Phi)/g_s}$$



Quantum Spectral Curves
(Riemann Surface + CFT)

Quantization

Geometry



*geometric
object*

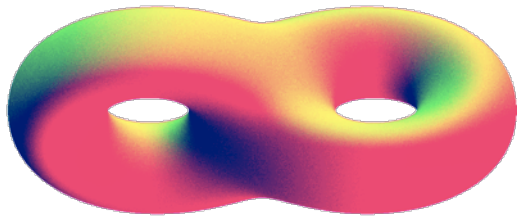
Algebra

$$Z(K) \in \mathbb{C}$$

*quantum
invariant*

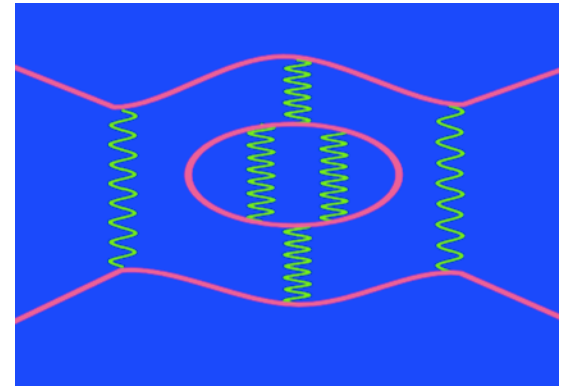
Emergence

Geometry



*effective
geometry*

Algebra



*quantum
system*

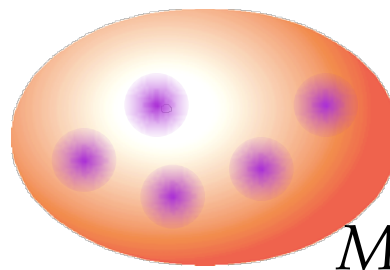
Supersymmetric Gauge Theory: N=4

Path integral localizes to instantons

$$Z = \int DA \cdot e^{-S} = \sum_{n \geq 0} d(n) q^n, \quad q = e^{2\pi i \tau}, \quad \tau \text{ coupling}$$

Contribution of moduli space

$$Mod_{N,n}(M)$$



M

$SL(2, \mathbf{Z})$ S-duality in N=4 gauge theories \leftrightarrow modular invariance of a quantum CFT on 2-torus

$$\tau \rightarrow \frac{a\tau + b}{c\tau + d}$$



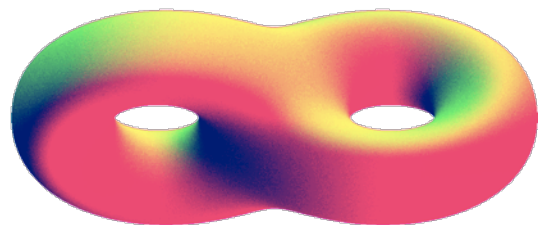
T^2

Related to 2d CFT characters

$$Z = \text{Tr } q^{L_0} = \chi(q)$$

Supersymmetric Gauge Theory: N=2

Seiberg-Witten curve



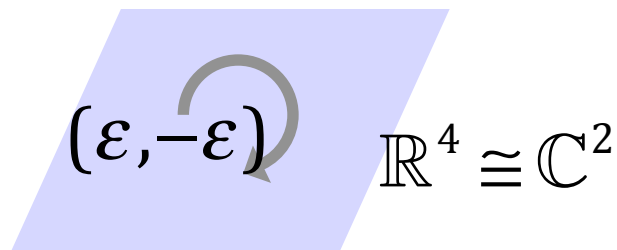
genus g

(Σ, ω) , $\omega =$ meromorphic one-form

Effective gauge coupling $\tau_{ij}(X)$, BPS masses $\oint_C \omega$

Nekrasov deformation

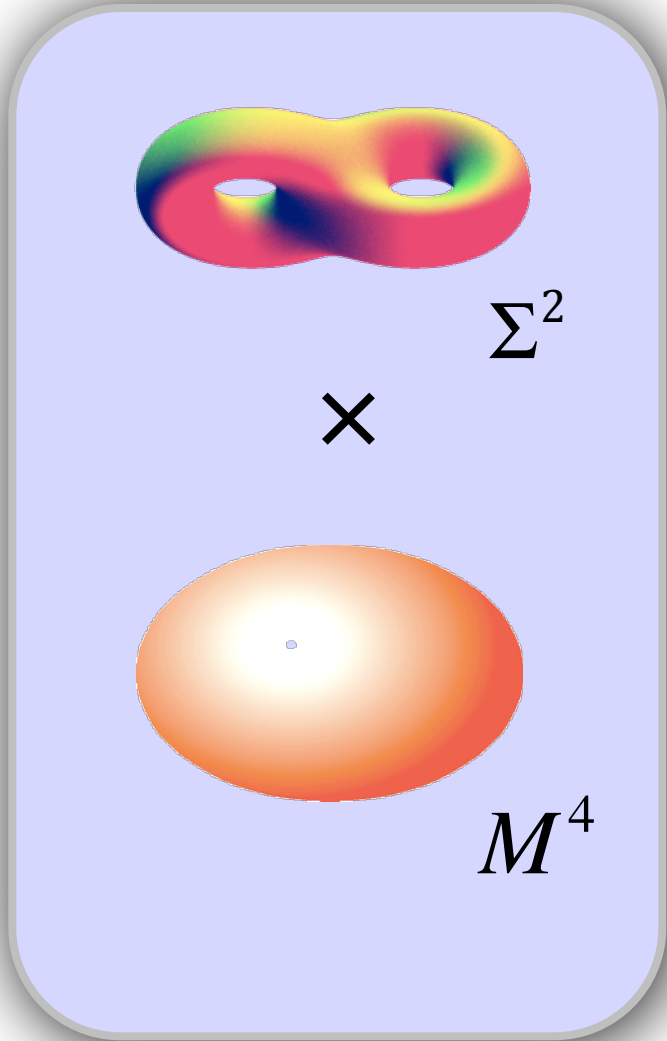
$$Z = e^F, \quad F = \sum_g \epsilon^{2g-2} F_g,$$



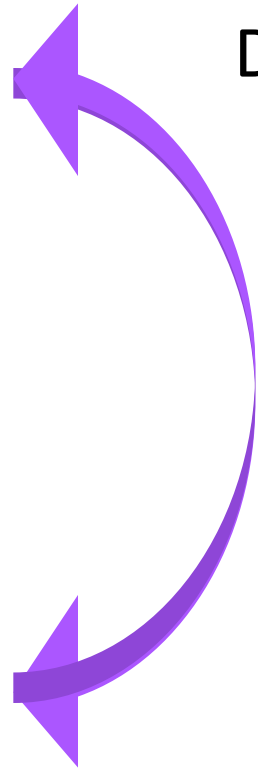
$$F_0 : \mu_i = \int_{A_i} \omega, \quad \frac{\partial F_0}{\partial \mu_i} = \int_{B_i} \omega$$

$$F_1 = \frac{1}{2} \log \det \Delta$$

D=6 (2,0) Tensor (2-form) SCFT



D=2 CFT

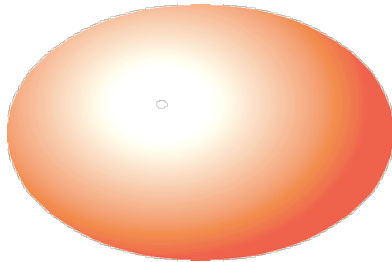


D=4 Gauge Theory

Geometry/Algebra Duality

Gauge Theories

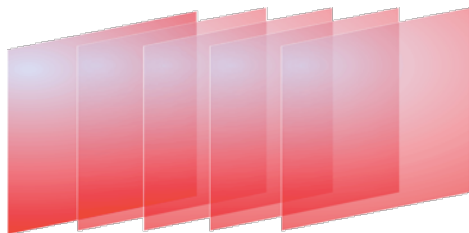
Geometry



M^4

4-manifold

Algebra



G

gauge group
(vector bundle)

Conformal Field Theory

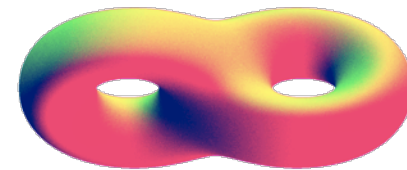
Algebra



\hat{g}

affine symmetries

Geometry



Σ^2

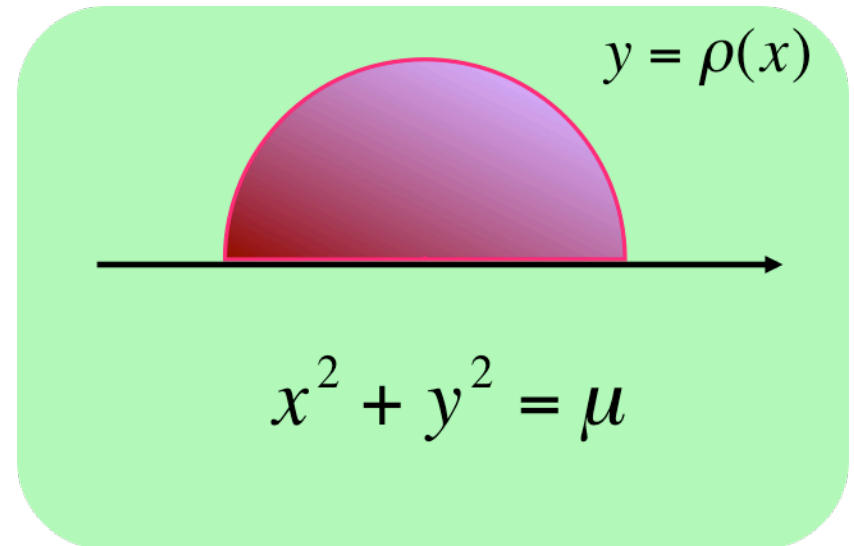
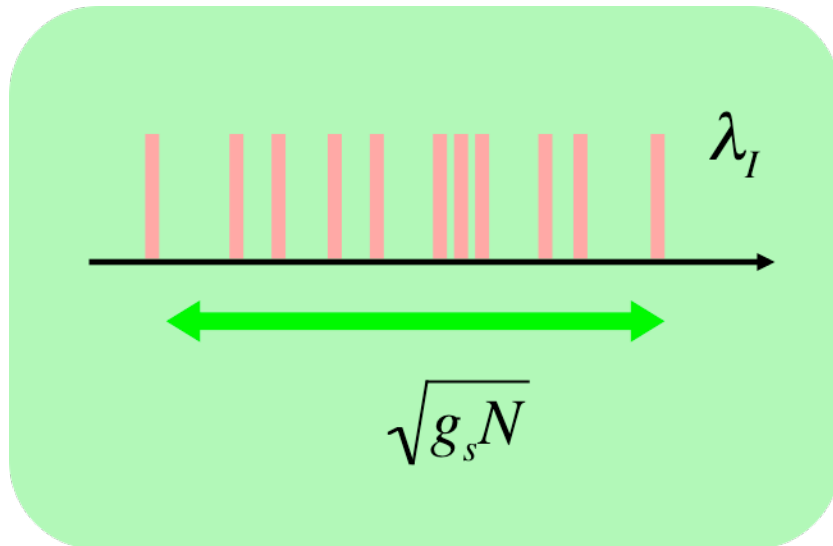
algebraic curve
Riemann surface

Wigner's Random Matrix Model

$$\lim_{N \rightarrow \infty} Z_N, \quad Z_N = \frac{1}{\text{Vol } U(N)} \int_{N \times N} d\Phi \cdot e^{-\text{Tr } \Phi^2 / g_s}$$

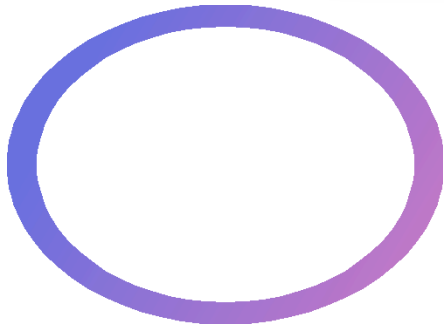
Eigenvalue distribution in 't Hooft limit

$$N \rightarrow \infty, \quad g_s \rightarrow 0, \quad Ng_s = \mu = \text{fixed}$$

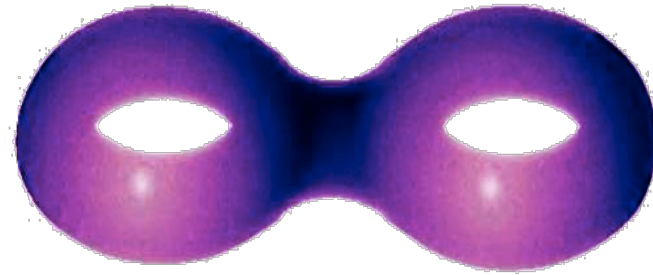


String Partition Function

$$Z = \exp \sum_{g \geq 0} g_s^{2g-2} F_g(\mu)$$



ribbon diagrams
1/N expansion



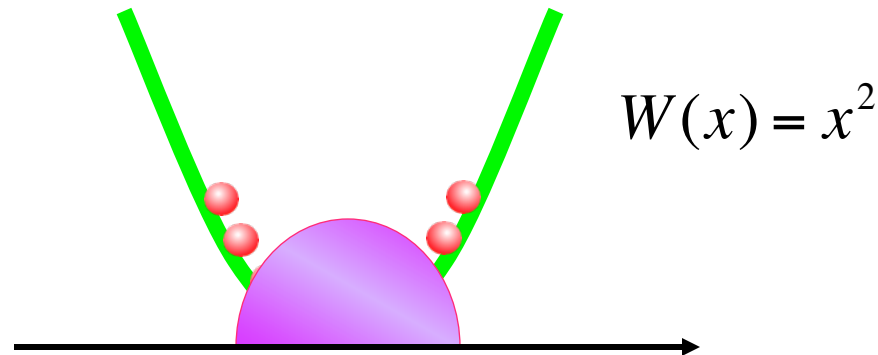
genus g
string surface

Eigenvalue Dynamics

$$Z_{matrix} = \int d^N \lambda \cdot \prod (\lambda_I - \lambda_J)^2 \cdot e^{-\sum_I \lambda_I^2 / g_s}$$

Effective action (repulsive Coulomb force)

$$S_{eff} = \sum_I \lambda_I^2 - 2g_s \sum_{I < J} \log(\lambda_I - \lambda_J)$$



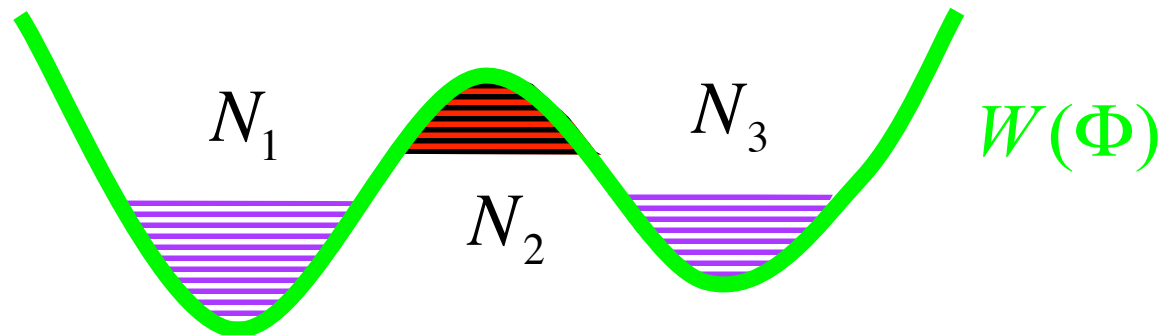
General Matrix Model

$$Z_{matrix} = \frac{1}{\text{Vol } U(N)} \int_{N \times N} d\Phi \cdot e^{\text{Tr } W(\Phi)/g_s}$$

't Hooft limit

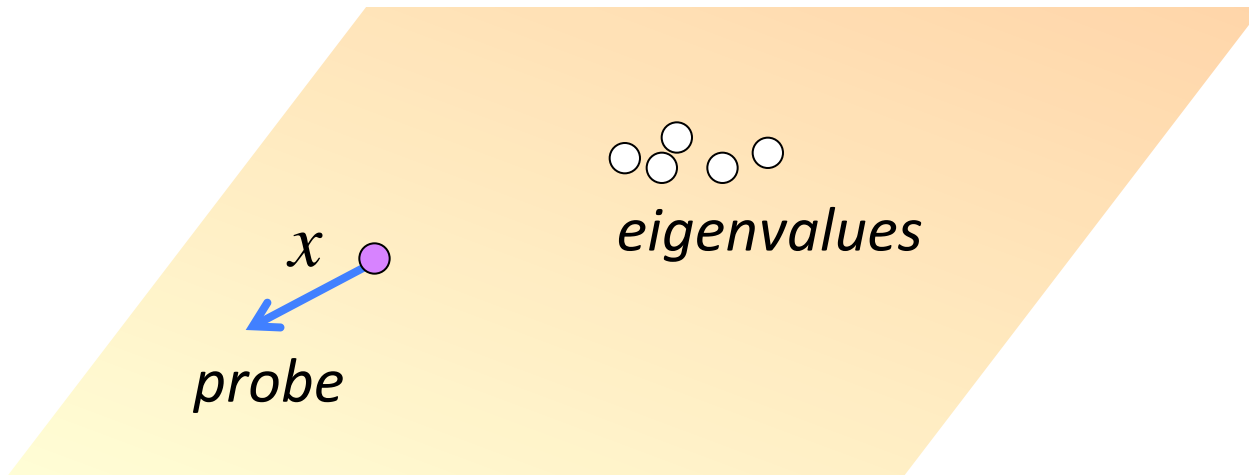
$$N_I \rightarrow \infty, \quad g_s \rightarrow 0, \quad N_I g_s = \mu_I = \text{fixed}$$

Filling fractions (perturbative expansion)



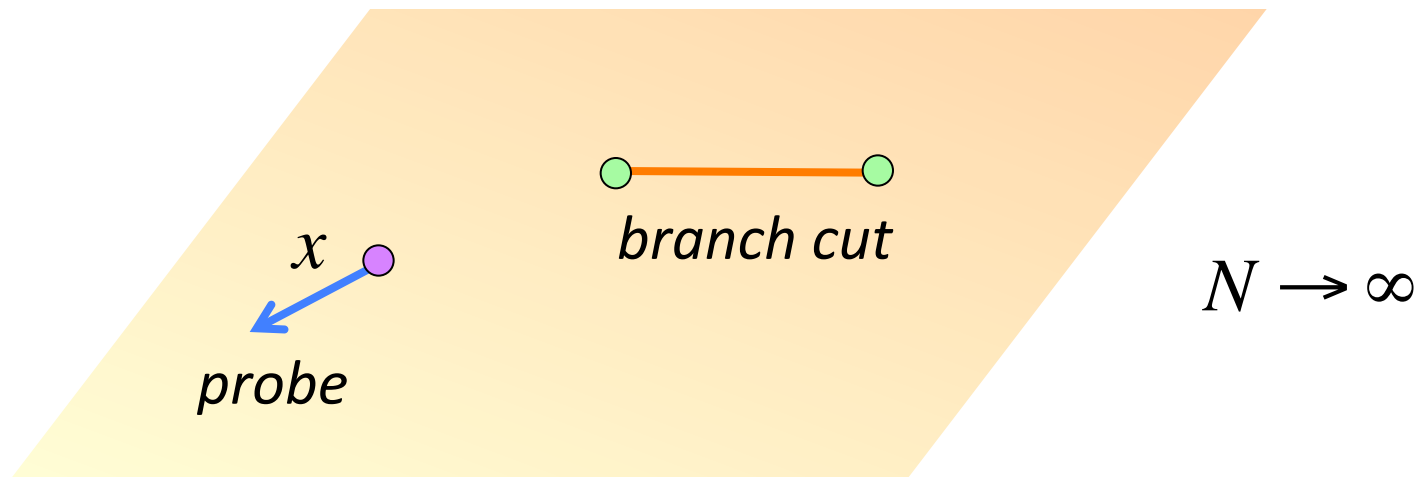
Resolvent

$$ydx = dS_{eff} = dW(x) + 2g_s \text{Tr} \frac{dx}{\Phi - x}$$



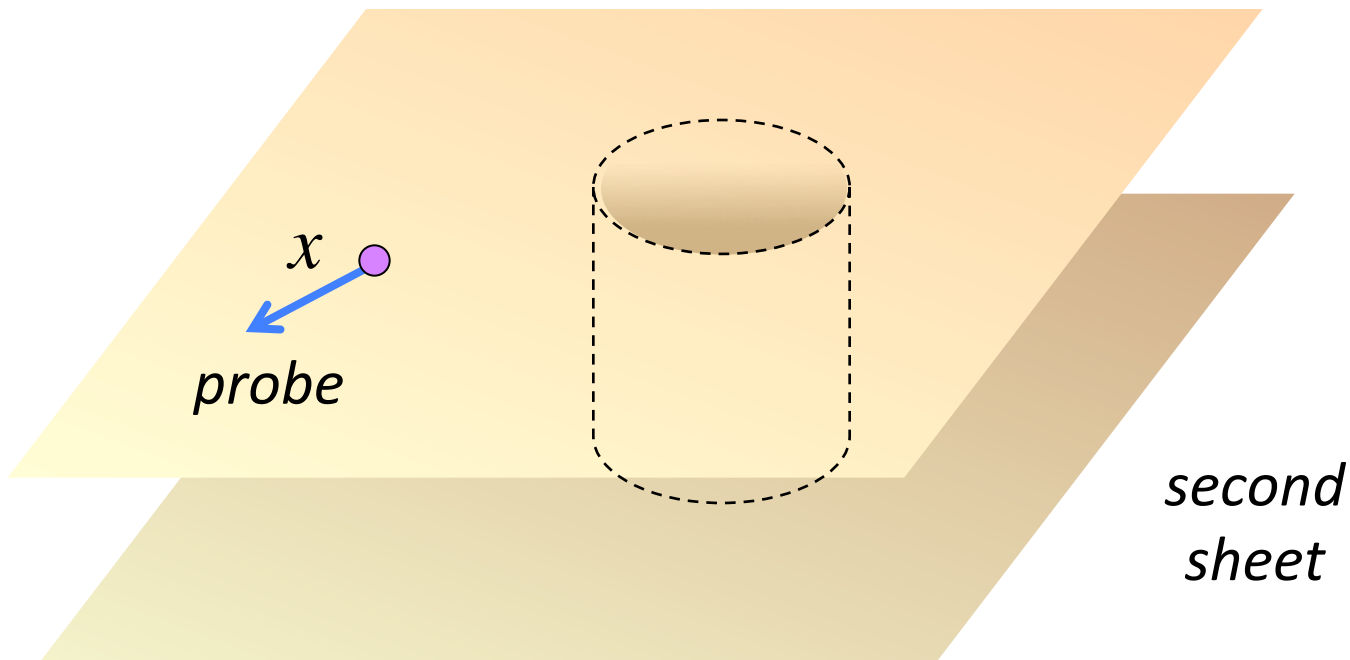
Resolvent

$$y dx = dS_{\text{eff}} = dW(x) + 2g_s \text{Tr} \frac{dx}{\Phi - x}$$



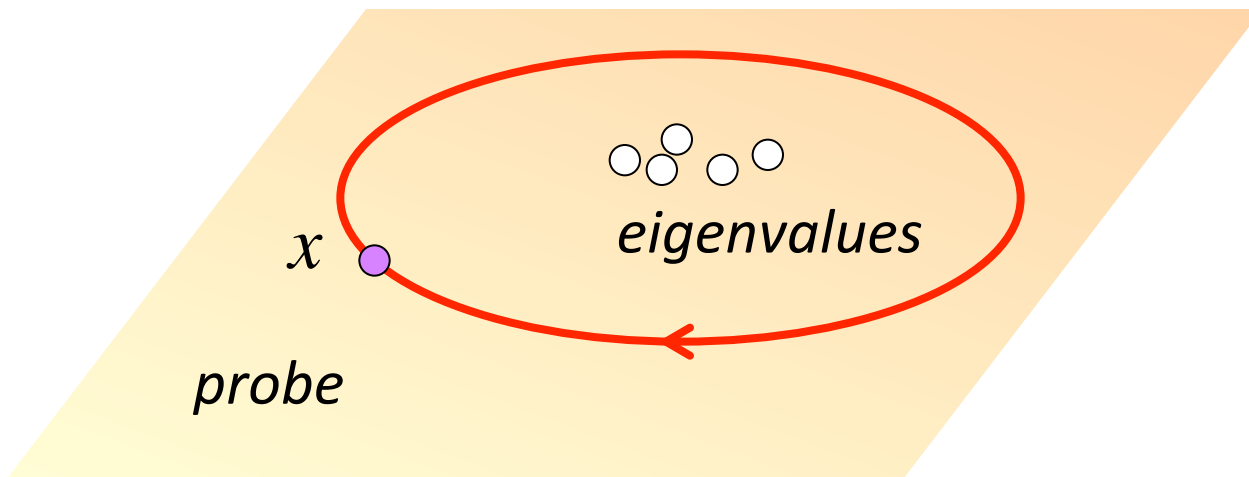
Effective Geometry

$$y^2 = W_n'(x)^2 + f_{n-1}(x) = P_{2n}(x)$$



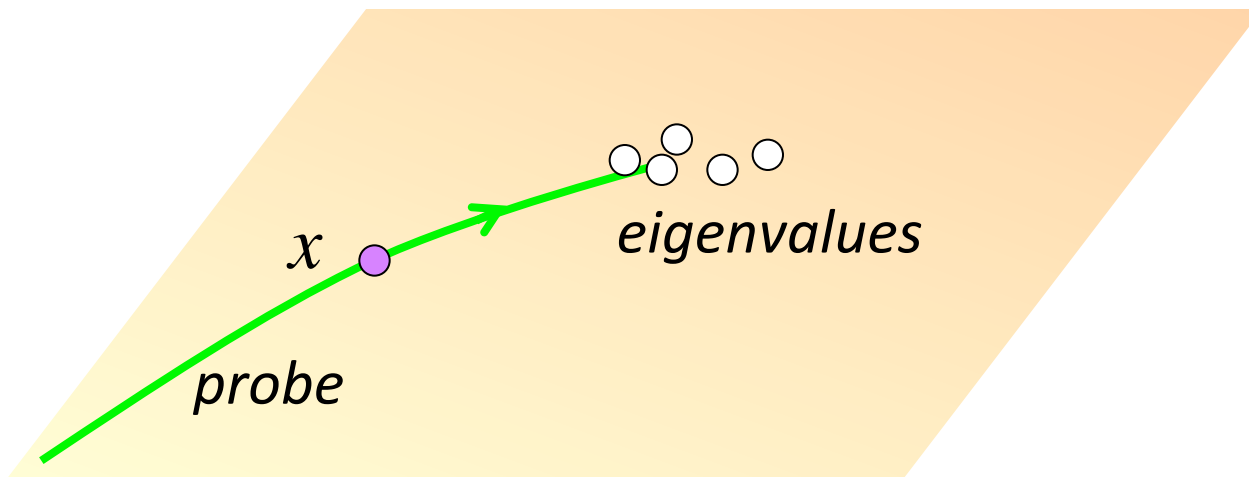
Periods (A)

$$\oint_{A_I} y dx = g_s N_I = \mu_I$$



Periods (B)

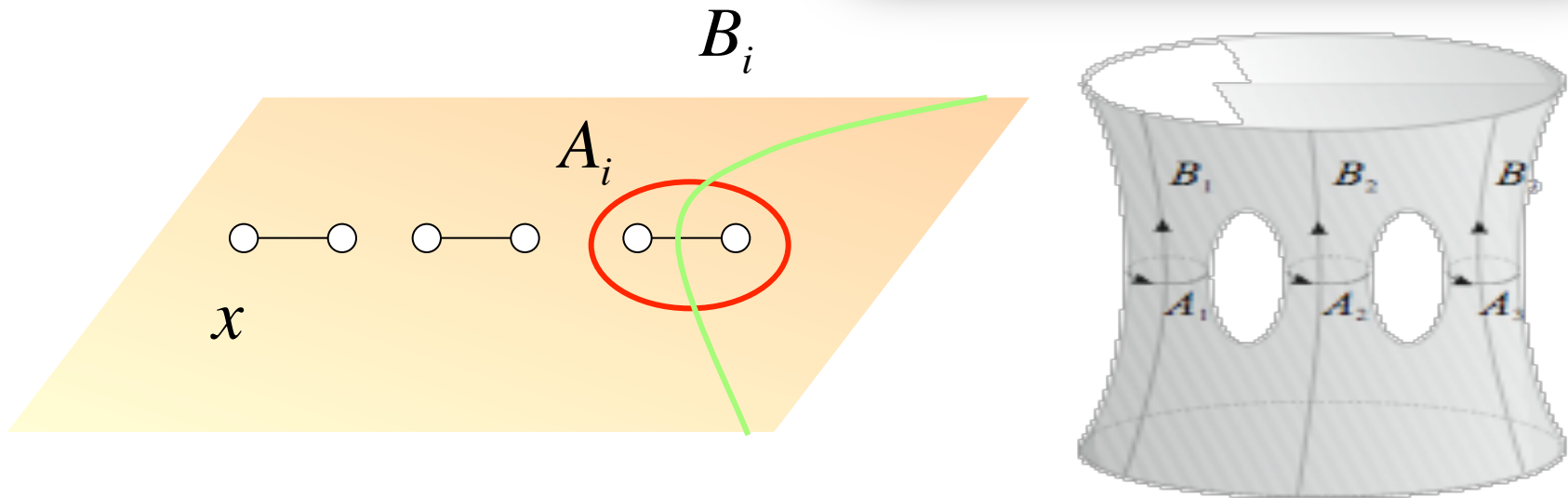
$$\int_{B_I} y dx = \frac{\partial F}{\partial \mu_I}$$



Spectral Curve

Hyperelliptic curve

$$\Sigma: y^2 = W'(x)^2 + f(x)$$



Quantum invariants of Σ , $\omega = ydx$

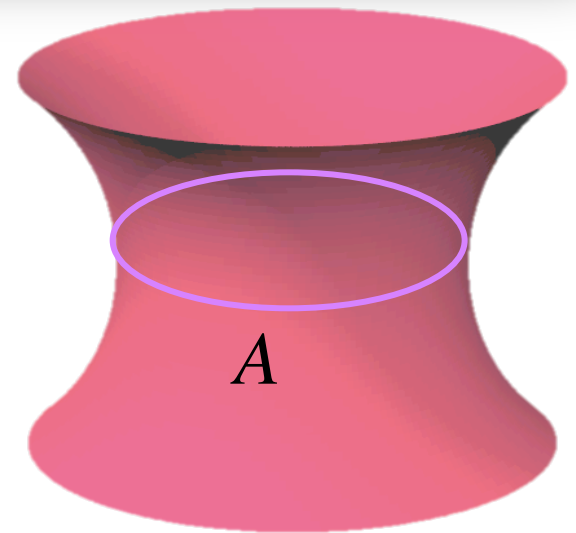
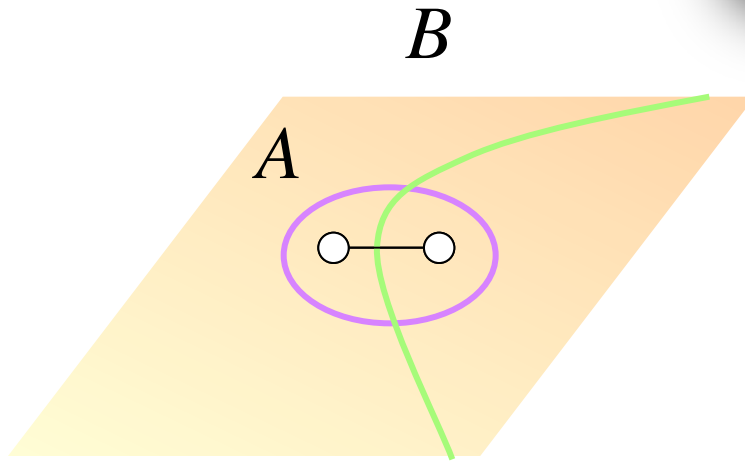
$$\mu_i = \oint_{A_i} ydx$$

$$\frac{\partial F_0}{\partial \mu_i} = \oint_{B_i} ydx$$

Wigner's Semi-Circle

Rational spectral curve

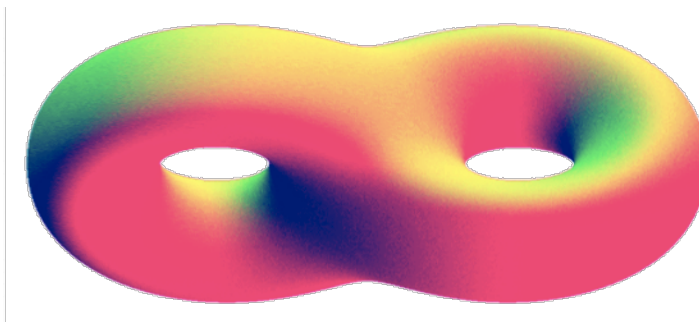
$$\Sigma: y^2 = x^2 + \mu$$



$$F_0 = \frac{1}{2} \mu^2 \log \mu = -\log \text{Vol}(U(N))$$

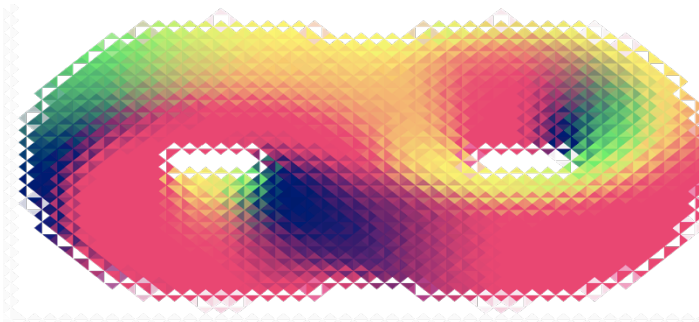
Quantum Curves

$N \rightarrow \infty$



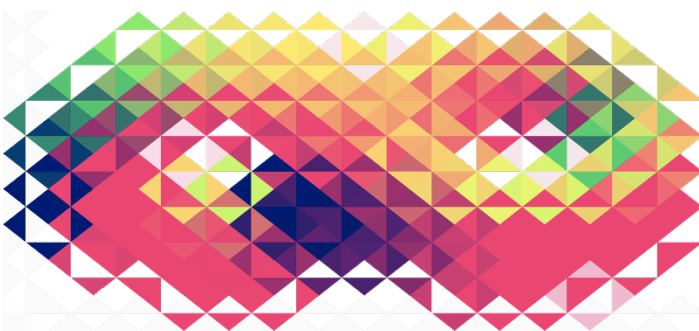
$$F_0(\mu)$$

$O(1/N)$



$$\sum_{g \geq 0} N^{2-2g} F_g(\mu)$$

N finite



non-perturbative

Complex Curves in Phase Space

algebraic curve = level set

$$\Sigma : H(x, y) = 0$$

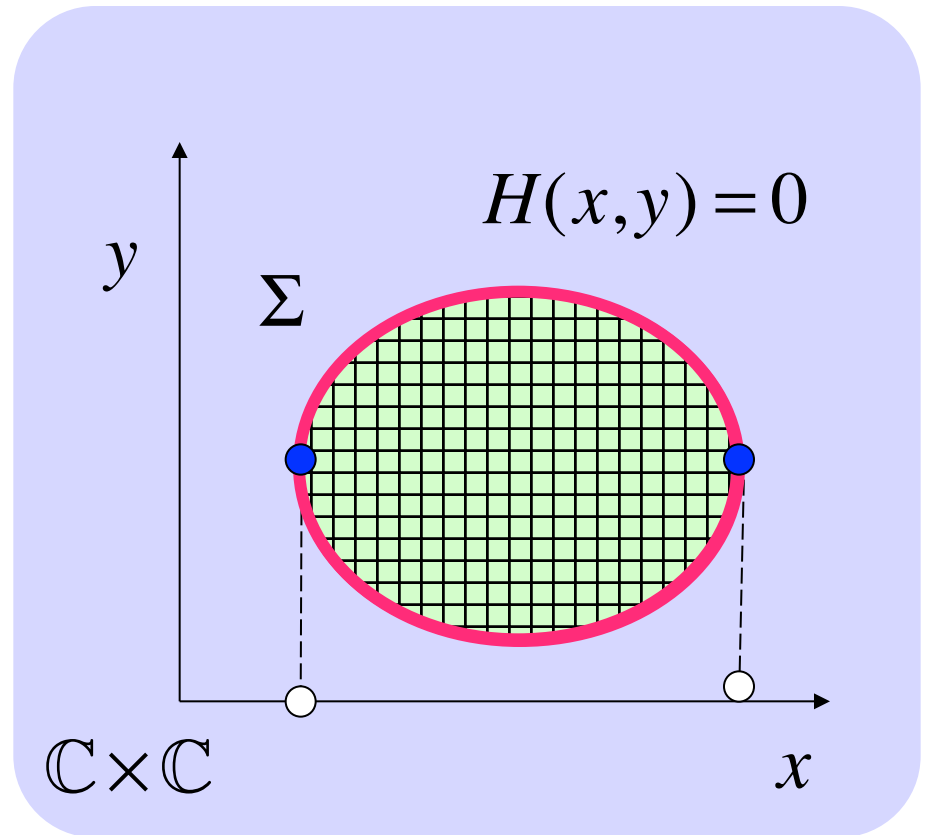
Hamilton-Jacobi theory

$$y = p(x)$$

branch pts = turning pts

Liouville form $\omega = ydx$

action
$$S(x) = \int_{x_*}^x \omega$$



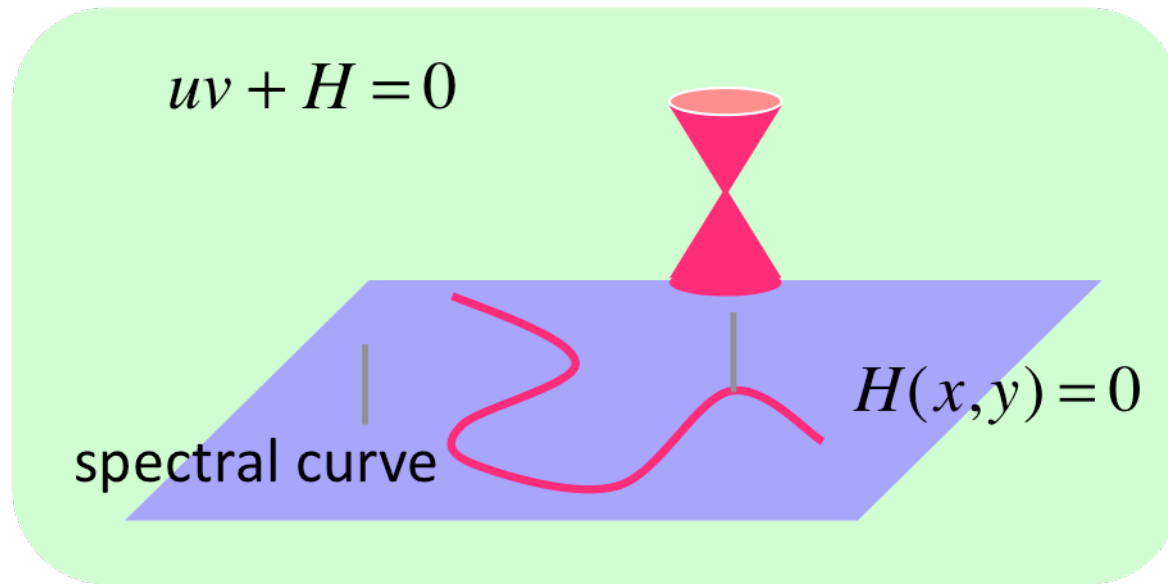
BS Quantization

$$dx \wedge dy \sim \hbar = g_s$$

Topological string on hypersurface CY_3

Calabi-Yau hypersurface in \mathbb{C}^4

$$X : uv + H(x, y) = 0$$



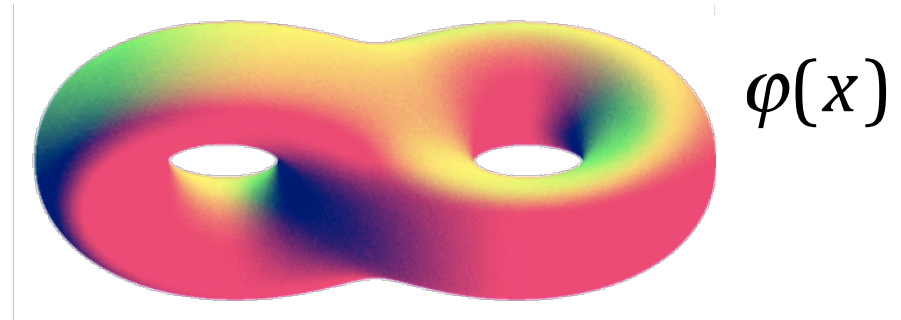
Period maps

$$\Omega = \frac{du}{u} \wedge dx \wedge dy \rightarrow \omega = ydx$$

Matrix models and chiral CFT

Eigenvalue density/matrix resolvent = collective boson field

$$\partial\varphi(x) = \text{Tr} \frac{1}{x - \Phi}$$



Free scalar field

$$Z_{matrix} = \int D\varphi \cdot e^{-S[\varphi]}, \quad S = \int (\partial\varphi)^2 + O(1/g_s)$$

Fermions = D-branes

$$\psi(x) = e^{i\varphi(x)} = \det(x - \Phi)$$
$$\psi^*(x) = e^{-i\varphi(x)} = \det(x - \Phi)^{-1}$$

Loop Equations & Virasoro Constraints

Invariance under diffeomorphism $x \rightarrow \tilde{x}(x)$

Generated by stress-tensor

$$T(x) = \frac{1}{2} (\partial\phi)^2(x)$$

Loop equations

$$\langle T(x) \rangle = W'(x)^2 + f(x)$$

SU(2) affine KM algebra structure

Currents

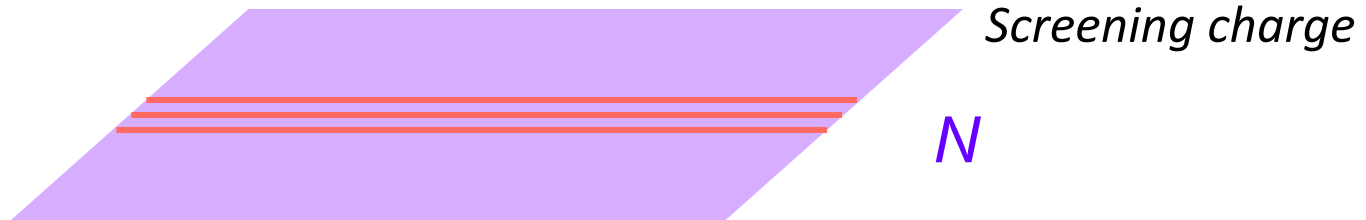
$$J_+ = e^{2i\varphi} = \psi_2^* \psi_1$$

$$J_3 = \partial\varphi = \psi_1^* \psi_1 - \psi_2^* \psi_2$$

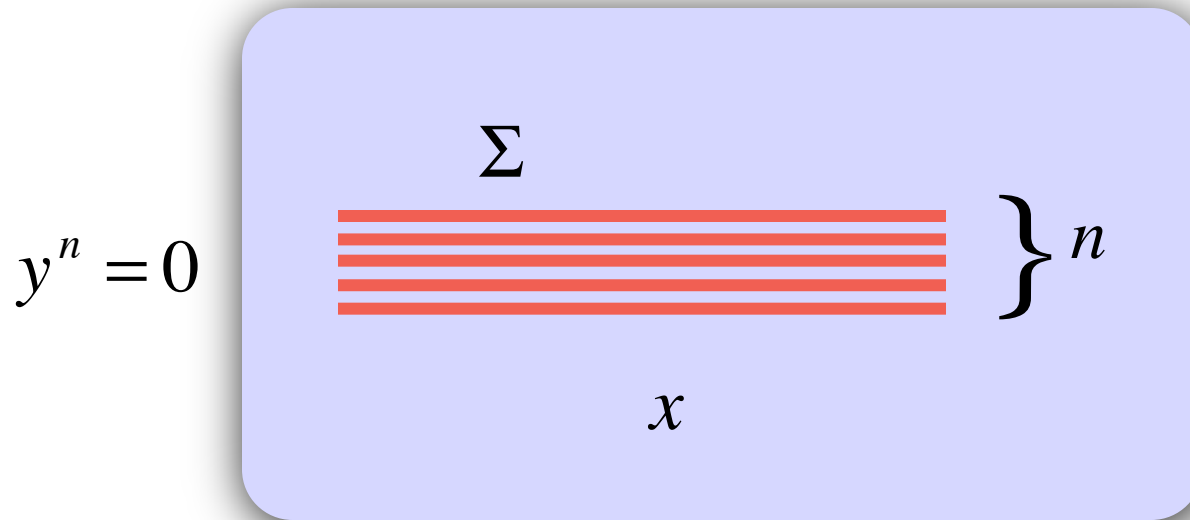
$$J_- = e^{-2i\varphi} = \psi_1^* \psi_2$$

In the limit $W=0$

$$Z_{matrix} = \left\langle \exp \int J_+ \right\rangle_N$$



Nilpotent Geometry



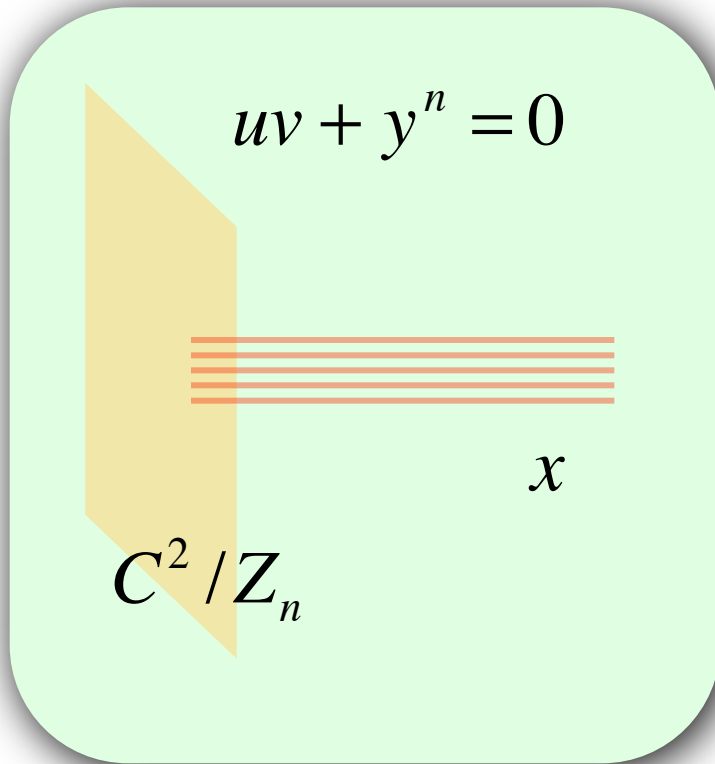
Nilpotent algebra

$$y^n = 0$$

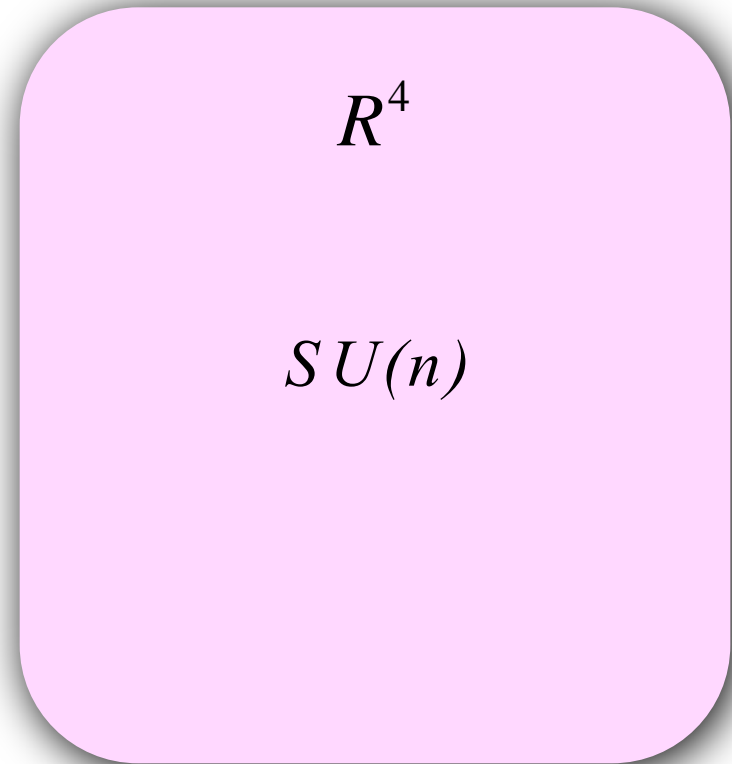
$$\hat{y} = g_s \partial_x + \begin{pmatrix} 0 & 0 & \dots & 0 \\ 1 & 0 & & \vdots \\ & \ddots & \ddots & 0 \\ 0 & & 1 & 0 \end{pmatrix}$$

Geometry of SU(n) Gauge Theory

A_{n-1} singularity



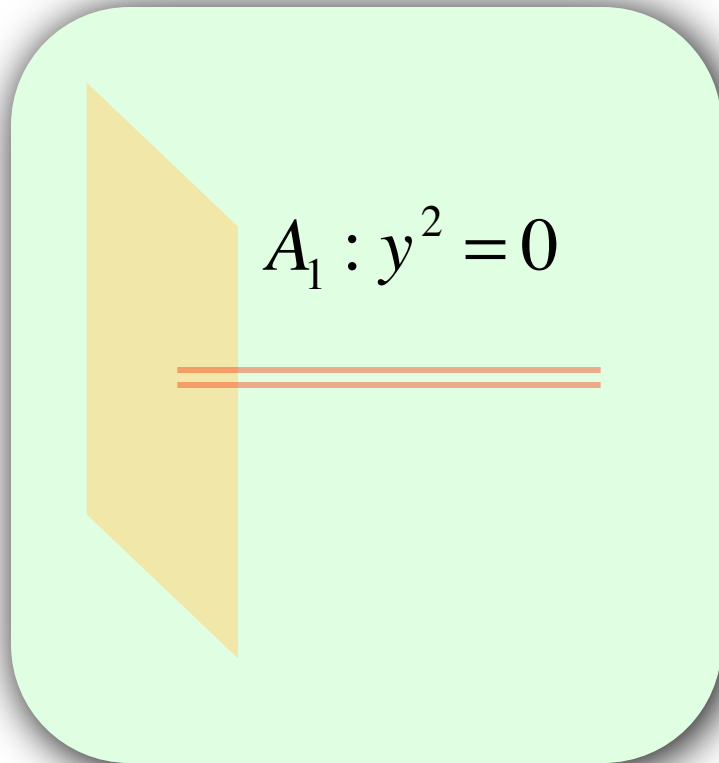
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Calabi-Yau
(internal space)

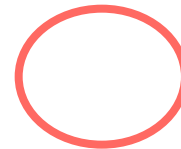
Space-Time

Singularities & Matrix Models



$$\int d\Phi \cdot e^{-a\text{Tr} \Phi^2}$$

$$\lim_{a \rightarrow 0} \{y^2 = a^2(x^2 + \mu)\}$$



Matrix Models & Liouville Theory

$$Z_{matrix} = \int d^N x \cdot \prod (x_I - x_J)^2 \cdot e^{\sum W(x_I)/g_s}$$

CFT Coulomb gas correlation function

$$Z_{matrix} = \left\langle \int d^N x \cdot \prod e^{2i\varphi(x_i)} \cdot e^{\oint \partial\varphi W/g_s} \right\rangle_N$$

Screening charge

$$Z_{matrix} = \left\langle e^{\int e^{2i\varphi}} \cdot e^{\oint \partial\varphi W/g_s} \right\rangle_N$$

Large N : Coulomb gas picture.

Beta Ensemble

General Liouville theory $b \neq i, \varepsilon_1 \neq -\varepsilon_2$

$$Z = \left\langle \exp \int e^{b\varphi} \right\rangle_N$$

Generalized matrix model measure

$$Z = \int d^N x \cdot \prod (x_I - x_J)^{2\beta} e^{\frac{\sqrt{\beta}}{g_s} \sum_I W(x_I)}, \quad \beta = -b^2$$

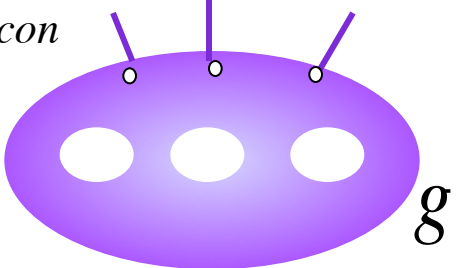
$$\beta = 1/2 : SO(N), \quad \beta = 2 : Sp(N)$$

Generalized loop equations

$$T(x) = \frac{1}{2} (\partial\varphi)^2 + Q \partial^2 \varphi$$

Eynard-Orantin: All-Genus Solution

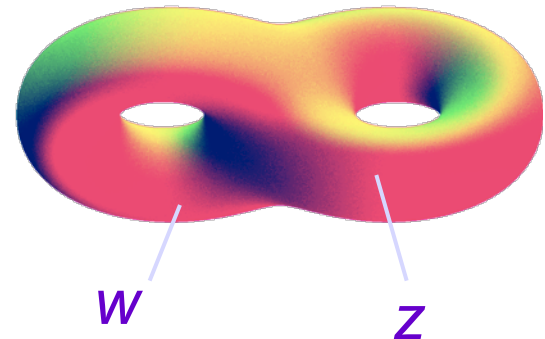
Recursion relations for correlation functions, completely geometric

$$\begin{aligned}
 W(x_1, \dots, x_n) &= \left\langle \text{Tr} \frac{1}{x_1 - \Phi} \dots \text{Tr} \frac{1}{x_n - \Phi} \right\rangle_{con} \\
 &= \langle \partial\varphi(x_1) \dots \partial\varphi(x_n) \rangle_{con} \\
 &= \sum_g g_s^{2g-2} W_g
 \end{aligned}$$


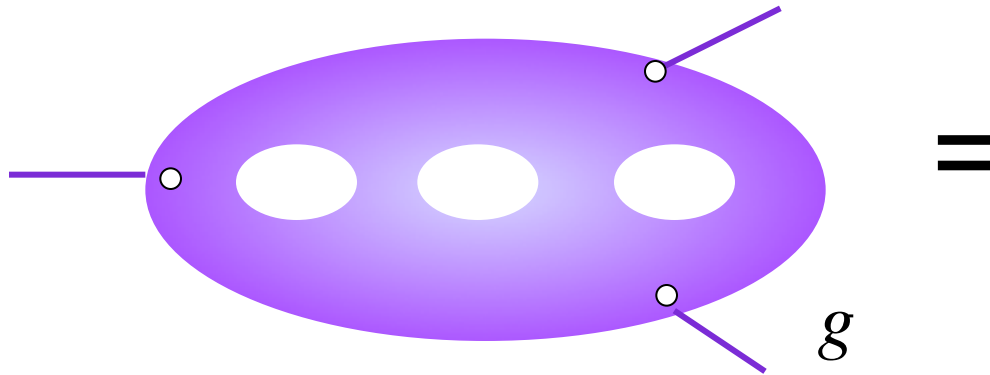
A purple genus g surface with three holes. Three small white circles are marked on the top boundary, each with a purple line segment extending upwards, representing insertion points for the correlation function.

Lowers genus, adds more insertions. Reduces to Bergmann kernel

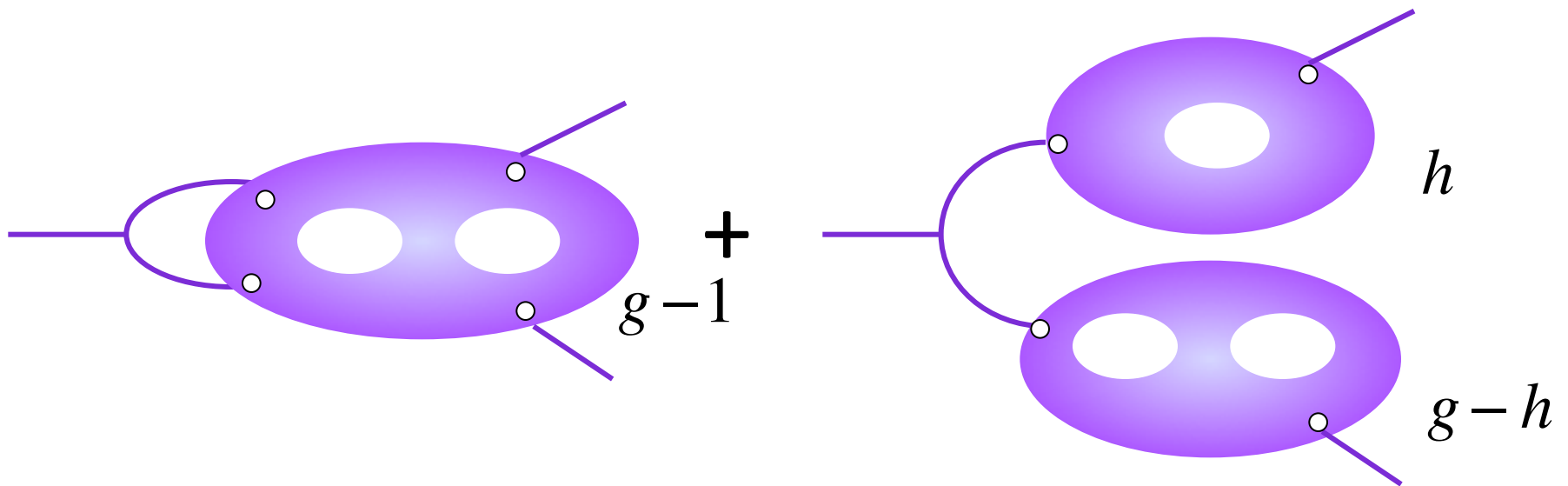
$$W_0(w, z) = B(w, z) dw dz$$



Recursion Relations

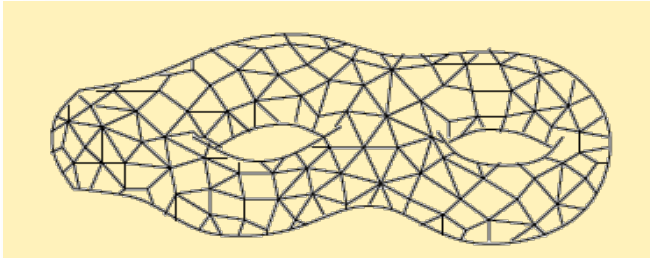


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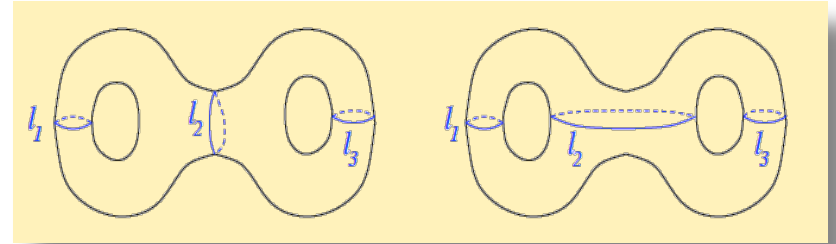
Quantum Spectral Curves

Random surfaces



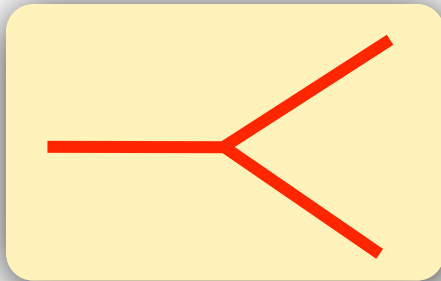
$$y^p = x^q$$

Geometry of \mathcal{M}_g



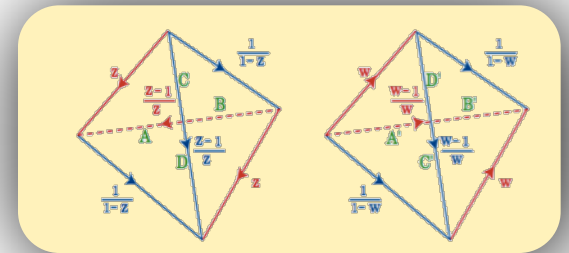
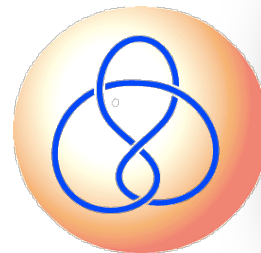
$$y = \sin \sqrt{x}$$

Toric Calabi-Yau



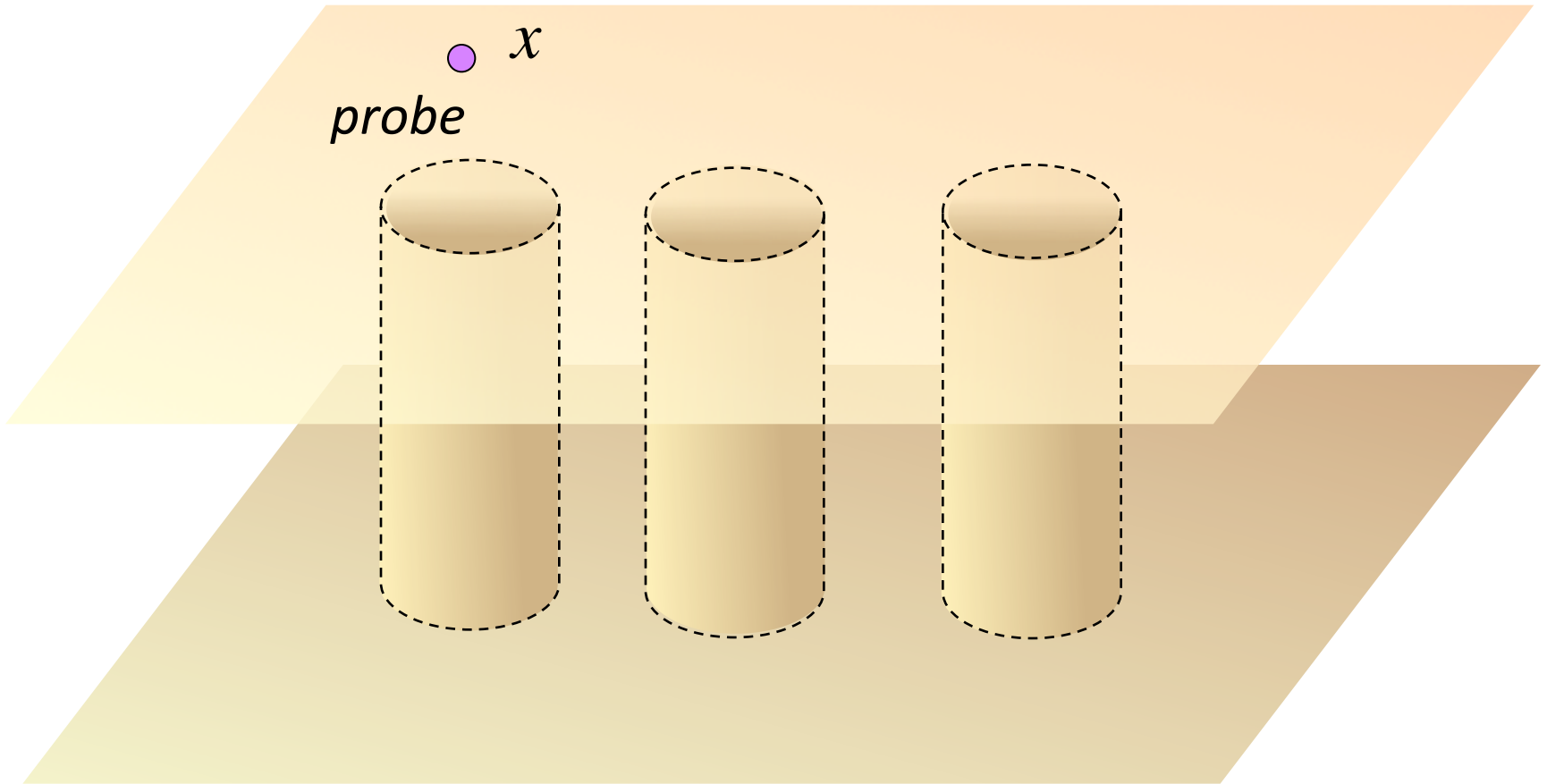
$$e^x + e^{-y} = 1$$

Hyperbolic knots



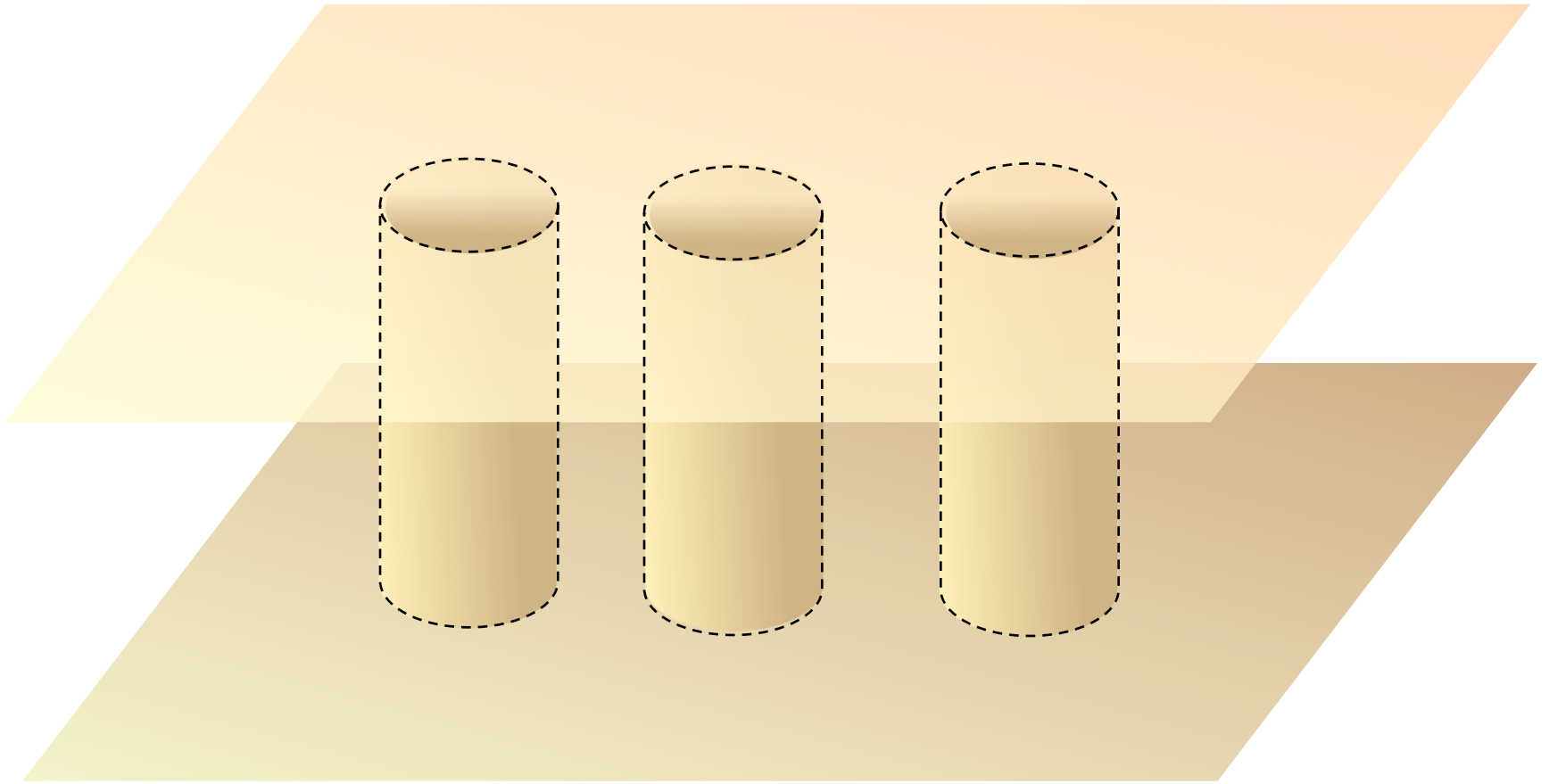
$$A(e^x, e^y) = 0$$

Effective Geometry

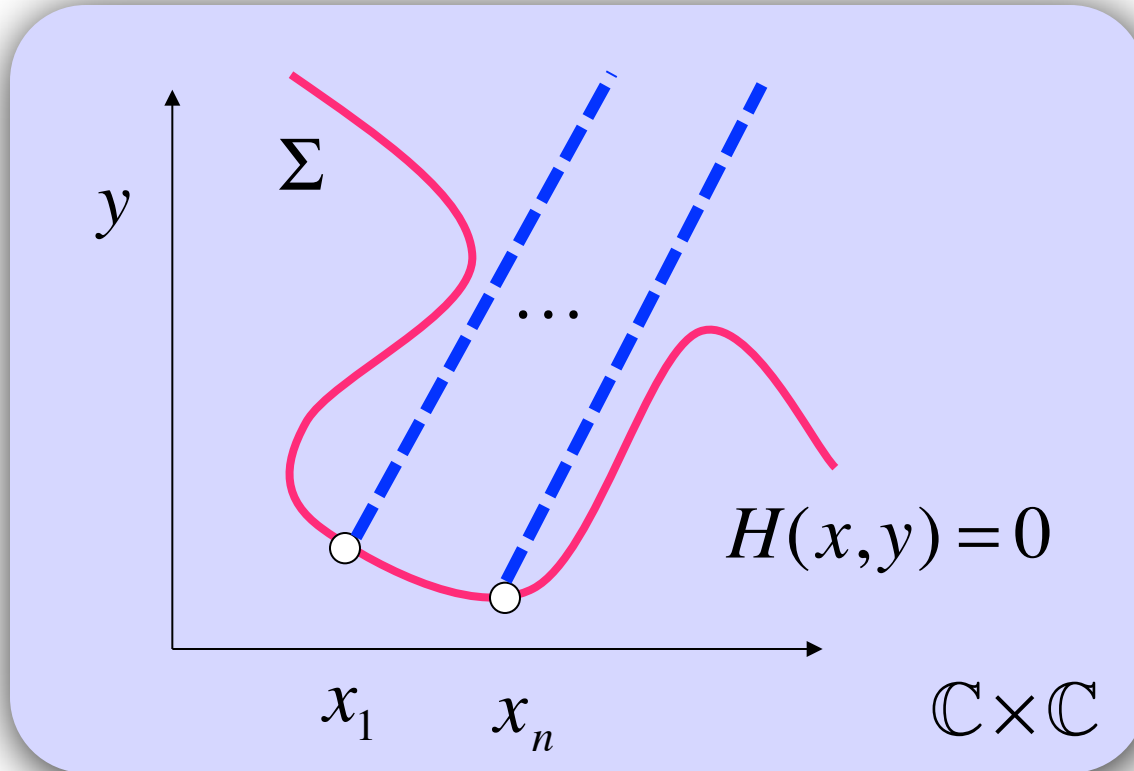


Are both sheets visible?

● x
probe



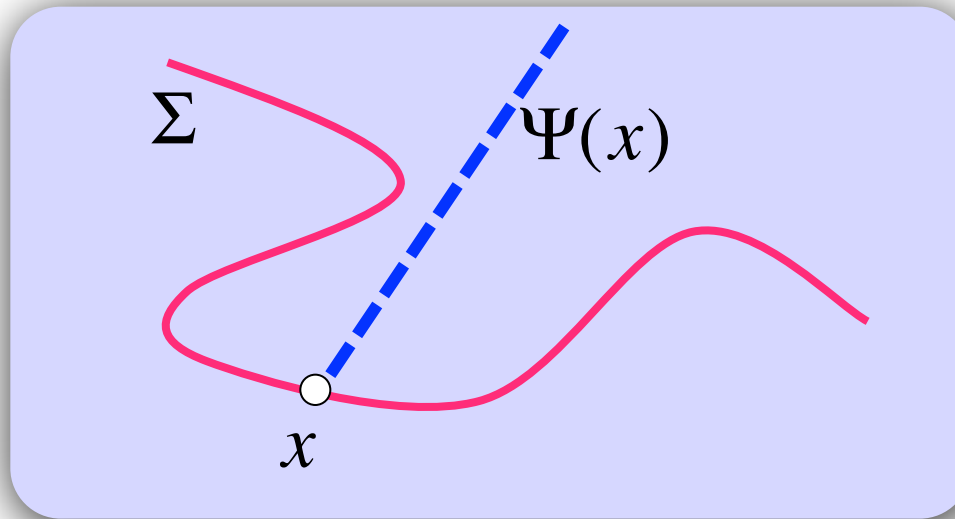
Topological D-Branes



Open string partition function

$$\Psi(x_1, \dots, x_n) = \langle \psi(x_1) \cdots \psi(x_n) \rangle = \left\langle \prod_i \det(\Phi - x_i) \right\rangle$$

Quantum Wave Function



Single brane wave function $\Psi(x)$

$$\hat{H}\Psi = 0, \quad \hat{H} = \hat{H}(\hat{x}, \hat{y})$$

Quantization of phase space

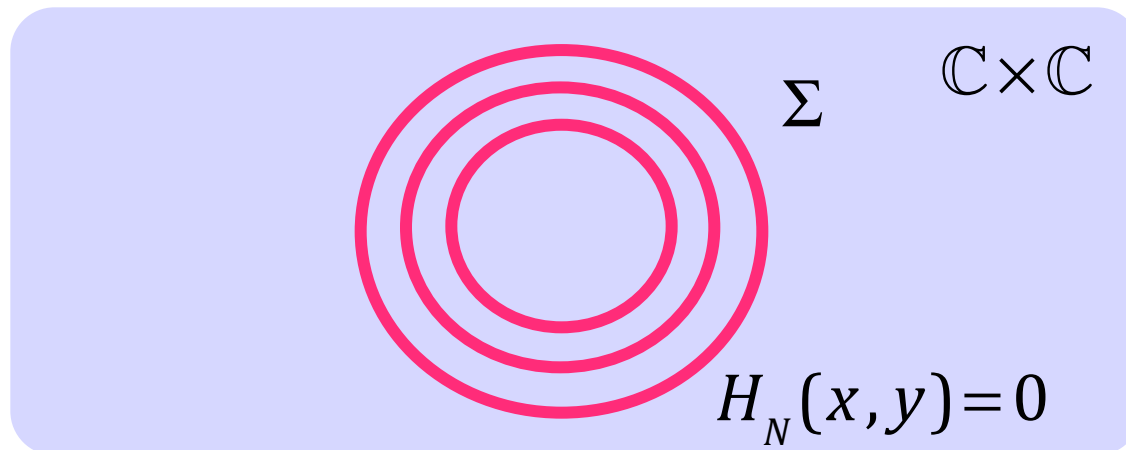
$$\hat{y} = -g_s \frac{\partial}{\partial x}, \quad [\hat{x}, \hat{y}] = g_s$$

Gaussian Matrix Model

$$\Psi_N(x) = \left\langle \det(\Phi - x) \right\rangle_N = H_{N-1}(x) \cdot e^{-x^2/2}$$

Eigenfunctions of harmonic oscillator

$$\left(-g_s^2 \frac{\partial^2}{\partial x^2} + x^2 - g_s(2N-1) \right) \Psi_N(x) = 0$$



Antibranes

$$\Psi_N^*(x) = \left\langle \frac{1}{\det(\Phi - x)} \right\rangle_N = \int \frac{\Psi_N(y)}{y - x} dy$$

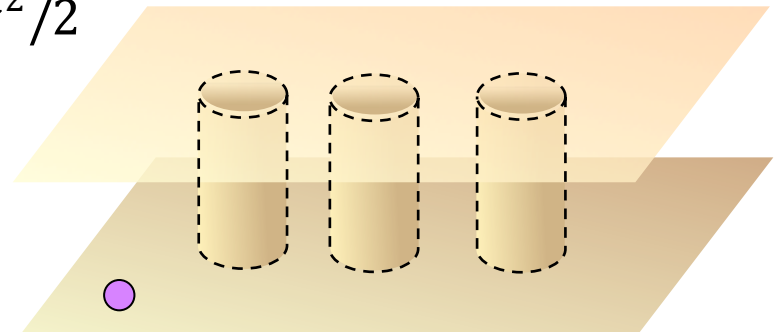
Non-normalizable eigenfunctions of harmonic oscillator

$$\hat{H}_N \Psi_N^*(x) = 0$$

WKB behavior

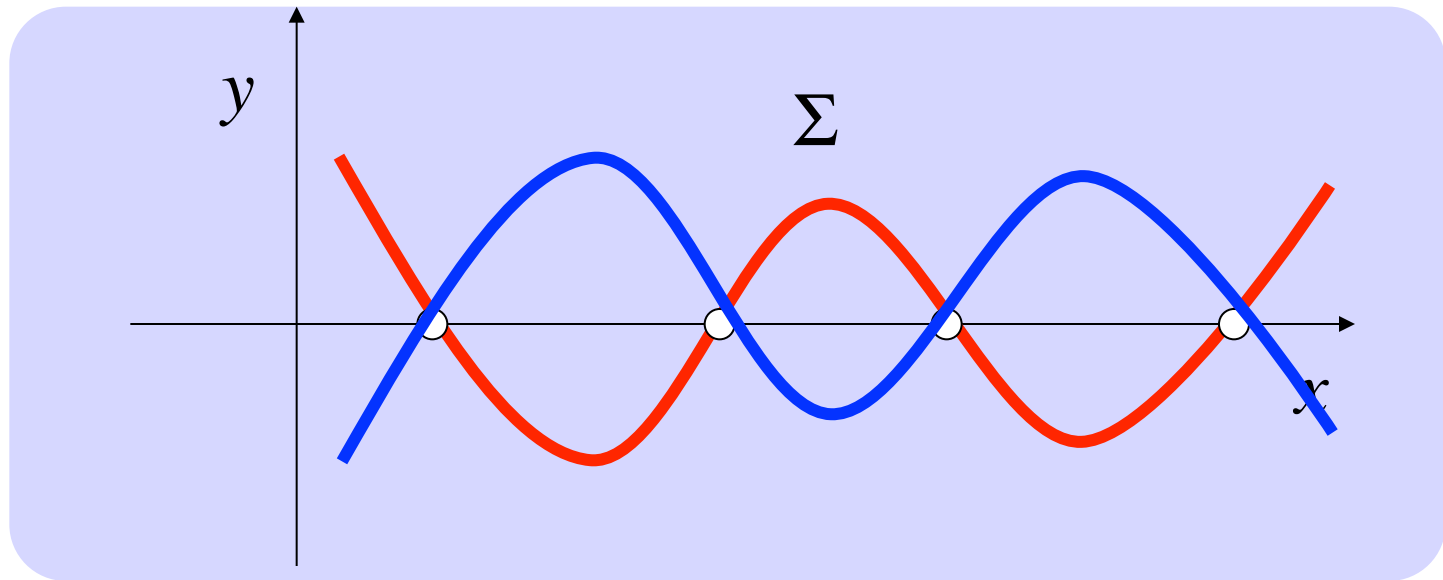
$$\Psi^*(x) \sim e^{+x^2/2}$$

Probe other sheet



Classical Spectral Curve

$$y^2 = W'(x)^2$$



branes: $y = W'(x)$

antibranes: $y = -W'(x)$

Super Gauge Theories

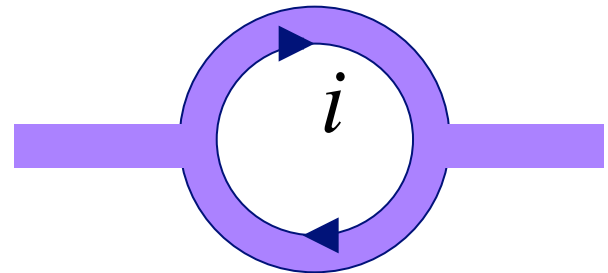
Super gauge group $U(N+k|k)$

Symmetries of $v = (z, \theta) \in \mathbb{C}^{N+k|k}$, $|v|^2 = \sum_{i=1}^{N+k} |z_i|^2 + \sum_{a=1}^k \theta_a \bar{\theta}_a$

$$\Phi = \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \quad A, D \text{ even, } B, C \text{ odd.}$$

Large N, M expansion

$$\text{Str}(1) = \sum_{i=1}^{N+k} 1 - \sum_{j=1}^k 1 = N$$



Perturbatively

$$U(N+k|k) \cong U(N)$$

Super Matrix Models

Supergroup

$$Z_{N+k|k} = \frac{1}{\text{Vol } U(N+k|k)} \int d\Phi \cdot e^{\text{Str } W(\Phi)/g_s}$$

Volume Lie supergroup vanishes: need regularization

$$\int d\theta \cdot 1 = 0$$

Diagonalize: eigenvalues & anti-eigenvalues

$$\Phi \sim \begin{pmatrix} x_i & 0 \\ 0 & y_j \end{pmatrix} \quad S = \sum_{i=1}^{N+k} W(x_i) - \sum_{j=1}^k W(y_j)$$

Super Jacobian

$$d\Phi \rightarrow d^{N+k}x d^k y \frac{\prod_{i<j} (x_i - x_j)^2 \prod_{a<b} (y_a - y_b)^2}{\prod_{i,a} (x_i - y_a)^2}$$

Eigenvalues of charge +1 and -1 (Coulomb gas model)

Resolvent

$$y dx \sim \text{Str} \frac{dx}{\Phi - x} = \sum_i \frac{dx}{x_i - x} - \sum_i \frac{dx}{y_a - x}$$

Same periods as U(N) model

Non-Perturbative Corrections

[Vafa] $U(n+k|k)$ has more Casimirs than $U(N)$, e.g.

$$\Delta = \prod_{i,a} (x_i - y_a)$$

General deformation

$$W = \sum_n t_n \text{Str } \Phi^n$$

For example, for $U(2|1)$

$$\Delta = -\frac{1}{2} \left(\text{Str } \Phi^2 - (\text{Str } \Phi)^2 \right) = 0 \text{ for } U(1)$$

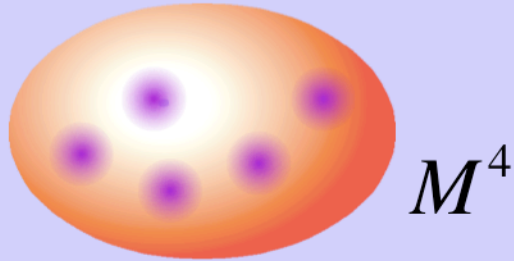
Non-Perturbative Corrections

So correlators $\langle \Delta \rangle_{N+k|k}$ are sensitive to k .

General pattern: calculation of partition function and correlation functions for $U(N+k|k)$ depend non-perturbatively (order e^{-N}) on k .

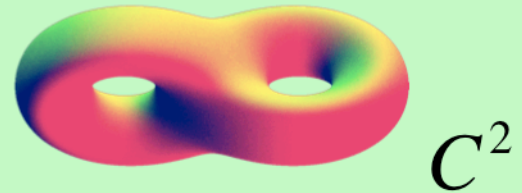
Working with supermatrix models seems to probe the full effective geometry.

Gauge Theories



self-dual connections
on 4-manifolds

Conformal Field Theory



2d QFT on algebraic curves
reps of infinite dim algebras

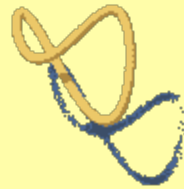
Matrix Models

$$\int_{N \times N} d\Phi \cdot e^{\text{Tr}W(\Phi)}$$



Eigenvalue distribution

(Topological) Strings



holomorphic curves,
mirror symmetry